

Close-to-contact dynamics of solids inside a viscous fluid

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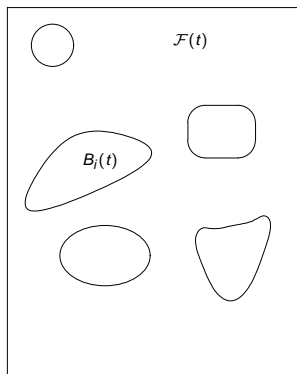


Figure 1. Rigid bodies in a fluid

Outline of the talks:

- 1 Classical equations
- 2 Study of the Cauchy problem in the viscous case
- 3 Close-to-contact motion of rigid bodies

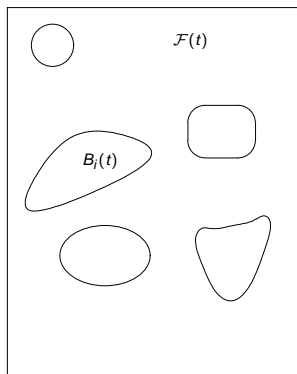


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Motivations :

- help numerical simulations
- find the limits of classical equations

I Classical equations

Solid description

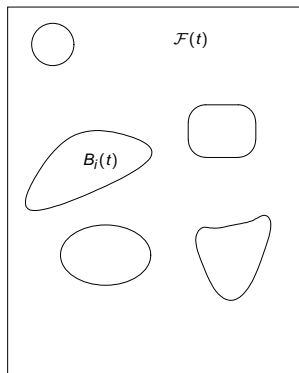


Figure 1. Rigid bodies in a fluid

Description of solid B_i

Shape $B_i \subset \mathbb{R}^n$ ($n = 2, 3$):

$$\partial B_i \in \mathcal{C}^{1,1}, \quad \int_{B_i} \mathbf{x} d\mathbf{x} = 0$$

Position $B_i(t)$:

$$B_i(t) = Q_i(t)B_i + \mathbf{G}_i(t),$$

where :

$$\mathbf{G}_i(t) \in \mathbb{R}^n, \quad Q_i(t)Q_i^\top(t) = \mathbb{I}_n.$$

$$Q_i(0) = \mathbb{I}_n$$

Solid description

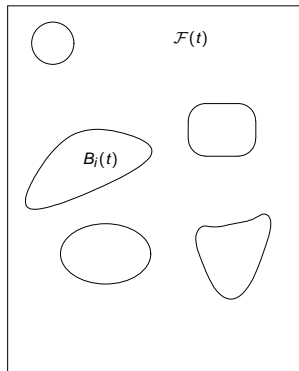


Figure 1. Rigid bodies in a fluid

Description of solid \mathcal{B}_i

Density $\rho_i \in (0, \infty)$

Mass $m_i \in (0, \infty)$,

$$m_i = \rho_i |B_i|$$

Inertia \mathbb{J}_i

$$\mathbf{a} \cdot \mathbb{J}_i \mathbf{b} = \rho_i \int_{B_i} [\mathbf{a} \times \mathbf{Q}_i \mathbf{y}] \cdot [\mathbf{b} \times \mathbf{Q}_i \mathbf{y}] \, d\mathbf{y}.$$

Solid dynamics

Solid kinematics : **Velocity-field**

$$\mathbf{u}_i = \mathbf{V}_i + \omega_i \times (\mathbf{x} - \mathbf{G}_i),$$

where $\mathbf{V}_i = \dot{\mathbf{G}}_i$, and $\omega_i \times " = " \dot{Q}_i Q_i^T$:

Solid dynamics.

$$(NL) \quad \left\{ \begin{array}{l} m_i \frac{d\mathbf{V}_i}{dt} = - \int_{\partial B_i(t)} \Sigma_i \mathbf{n}_i d\sigma_i \\ \frac{d[\mathbb{J}_i \omega_i]}{dt} = - \int_{\partial B_i(t)} [\mathbf{x} - \mathbf{G}_i] \times [\Sigma_i \mathbf{n}_i] d\sigma_i \end{array} \right.$$

Notations : Σ_i stress tensor, \mathbf{n}_i unit normal (toward $B_i(t)$).

Solid dynamics

Solid kinematics : Velocity-field

$$\mathbf{u}_i = \mathbf{V}_i + \omega_i \times (\mathbf{x} - \mathbf{G}_i),$$

where $\mathbf{V}_i = \dot{\mathbf{G}}_i$, and $\omega_i \times \mathbf{x} = \dot{\mathbf{Q}}_i \mathbf{Q}_i^T \mathbf{x}$:

$$\omega \times \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \omega \begin{pmatrix} -x_2 \\ x_1 \end{pmatrix}$$

Solid dynamics.

$$(NL) \quad \begin{cases} m_i \frac{d\mathbf{V}_i}{dt} = - \int_{\partial B_i(t)} \Sigma_i \mathbf{n}_i d\sigma_i \\ \frac{d[\mathbb{J}_i \omega_i]}{dt} = - \int_{\partial B_i(t)} [\mathbf{x} - \mathbf{G}_i] \times [\Sigma_i \mathbf{n}_i] d\sigma_i \end{cases}$$

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Solid dynamics.

$$(NL) \quad \begin{cases} m_i \frac{d\mathbf{V}_i}{dt} = - \int_{\partial B_i(t)} \Sigma_i \mathbf{n}_i d\sigma_i + \int_{B_i(t)} \rho_i \mathbf{f} \\ \frac{d[\mathbb{J}_i \omega_i]}{dt} = - \int_{\partial B_i(t)} [\mathbf{x} - \mathbf{G}_i] \times [\Sigma_i \mathbf{n}_i] d\sigma_i + \int_{B_i(t)} \rho_i [\mathbf{x} - \mathbf{G}_i] \times \mathbf{f} \end{cases}$$

Notations : Σ_i stress tensor, \mathbf{n}_i unit normal (toward $B_i(t)$).

Solid dynamics

Solid kinematics : Velocity-field

$$\mathbf{u}_i = \mathbf{V}_i + \omega_i \times (\mathbf{x} - \mathbf{G}_i),$$

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Notations : Σ_i stress tensor, \mathbf{n}_i unit normal (toward $B_i(t)$).

$$\begin{pmatrix} x_1 \\ x_2 \end{pmatrix} \times \begin{pmatrix} y_1 \\ y_2 \end{pmatrix} = x_1 y_2 - y_2 x_1$$

Fluid equations

Fluid description

- density ρ , viscosities μ, λ
- velocity-field \mathbf{u} , pressure p

Fluid equations Compressible Navier Stokes equations

$$(\text{CNS}) \quad \begin{cases} \rho(\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u}) = \operatorname{div} \Sigma + \rho \mathbf{f} \\ \partial_t \rho + \operatorname{div}(\rho \mathbf{u}) = 0 \\ p = a\rho^\gamma \end{cases} \quad \text{on } \mathcal{F}(t)$$

Fluid assumptions

- **Compressible** or Incompressible (constant-density)
- Inviscid or **Viscous (newtonian)**

$$\Sigma = \mu[\nabla \mathbf{u} + \nabla^\top \mathbf{u}] - (p + \lambda \operatorname{div} \mathbf{u}) \mathbb{I}_n$$

Fluid equations

Fluid description

- density ρ , viscosities μ, λ
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Fluid equations Incompressible Navier Stokes equations

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Fluid assumptions

- Compressible or **Incompressible** (constant-density)
- Inviscid or **Viscous (newtonian)**

$$\Sigma = \mu[\nabla \mathbf{u} + \nabla^T \mathbf{u}] - p \mathbb{I}_n$$

Fluid equations

Fluid description

- density ρ , viscosities μ, λ
- velocity-field \mathbf{u} , pressure p

Fluid equations Euler equations

$$(E) \quad \begin{cases} \rho(\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u}) & = \operatorname{div} \Sigma + \rho \mathbf{f} \\ \operatorname{div} \mathbf{u} & = 0 \\ \rho & = c_{st} \end{cases} \quad \text{on } \mathcal{F}(t)$$

Fluid assumptions

- Compressible or **Incompressible** (constant-density)
- **Inviscid** or Viscous Newtonian

$$\Sigma = -p \mathbb{I}_n$$

Interactions

- The fluid domain is fixed by solid positions :

$$\mathcal{F}(t) = \Omega \setminus \overline{\bigcup_i B_i(t)}.$$

- Boundary conditions are fixed by solid dynamics:

- Impermeability conditions :

$$(BC_i) \quad (\mathbf{u} - \mathbf{u}_i) \cdot \mathbf{n}_i = 0 \text{ on } \partial B_i(t), \quad \mathbf{u} \cdot \mathbf{n}_0 = 0 \quad \text{on } \partial\Omega$$

- No-slip conditions :

$$(BC_{ns}) \quad (\mathbf{u} - \mathbf{u}_i) \times \mathbf{n}_i = 0 \text{ on } \partial B_i(t), \quad \mathbf{u} \times \mathbf{n}_0 = 0 \quad \text{on } \partial\Omega$$

- Continuity of stress tensor :

$$\Sigma = \Sigma_i \quad \forall i.$$

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- Slip conditions :

$$(BC_s) \quad (\mathbf{u} - \mathbf{u}_i) \times \mathbf{n}_i = -\beta_i [\Sigma \mathbf{n}_i] \times \mathbf{n}_i \text{ on } \partial B_i(t), \quad \mathbf{u} \times \mathbf{n}_0 = -\beta_0 [\Sigma \mathbf{n}_0] \times \mathbf{n}_0 \quad \text{on } \partial\Omega$$

- Continuity of stress tensor :

$$\Sigma = \Sigma_i \quad \forall i.$$

Bibliography - Analytical point of view

- Incompressible inviscid fluid with impermeability conditions
 - **ideal fluids** : J.G. Houot & A. Munier '08
 - **general case** : J. Ortega *et al* '05, C. Rosier & L. Rosier '09, J.G. Houot *et al.* '11.
- Incompressible viscous fluid with no-slip boundary conditions
 - **Classical solutions** : C. Grandmont & Y. Maday '98, B. Desjardins & M.J. Esteban 99', M. Tucsnak & T. Takahashi '04, T. Takahashi '03, A.L. Silvestre & G.P. Gadi '05, P. Cumsille & T. Takahashi '08,
 - **Weak solutions** : C. Conca *et al* '00, M. Gunzburger *et al.* '00, B. Desjardins & M.J. Esteban 00', J. A. San Martin *et al.* '02, E. Feireisl '03
 - **Qualitative properties of solutions** : V.N. Starovoitov '02,'03, T.I. Hesla, '05, M.H. '07, M.H. & T. Takahashi & M.H. 09, D. Gérard-Varet & M. H. '10.
- Compressible viscous fluid with no-slip boundary conditions : B. Desjardins & M.J. Esteban 00', E. Feireisl '03,
- Incompressible viscous fluid with slippage ...

II Construction of solutions

Full system (v-FSIS)

$$N = 1 \quad \Omega = \mathbb{R}^2 \quad \mathbf{f} = 0 \quad \rho = 1$$

$$\begin{cases} (\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u}) &= \operatorname{div} \Sigma, \\ \operatorname{div} \mathbf{u} &= 0, \end{cases} \quad \forall (t, \mathbf{x}) \in \mathcal{Q}_T := \{(t, \mathbf{x}), \mathbf{x} \in \mathbb{R}^2 \setminus \bar{B}_1(t)\},$$

$$\mathbf{u}(t, \mathbf{x}) = \mathbf{V}_1(t) + \omega_1(t) \times (\mathbf{x} - \mathbf{G}_1), \quad \forall \mathbf{x} \in \partial B_1(t), \quad \lim_{|\mathbf{x}| \rightarrow \infty} \mathbf{u}(t, \mathbf{x}) = 0.$$

$$\begin{cases} m_1 \frac{d\mathbf{V}_1}{dt} &= - \int_{\partial B_1(t)} \Sigma \mathbf{n}_1 \, d\sigma_1, \\ J_1 \frac{d[\omega_1]}{dt} &= - \int_{\partial B_1(t)} [\mathbf{x} - \mathbf{G}_1] \times [\Sigma \mathbf{n}_1] \, d\sigma_1, \end{cases}$$

$$\Sigma = \Sigma(\mathbf{u}, p) := \mu(\nabla \mathbf{u} + \nabla^T \mathbf{u}) - p \mathbb{I}_2,$$

$$\mathbf{u}(0, \cdot) = \mathbf{u}^0, \quad \mathbf{G}_1(0) = \mathbf{G}_1^0, \quad \mathbf{V}_1(0) = \mathbf{V}_1^0, \quad \omega_1(0) = \omega_1^0, \quad Q_1(0) = \mathbb{I}_2.$$

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Change of geometry

T. Takahashi '03, M. Tucsnak & T. Takahashi '04

$t \mapsto (\mathbf{V}_1, \omega_1)$ is given

Local velocity around the solid body :

$$\mathbf{u}_1(t, \mathbf{x}) = \mathbf{V}_1(t) + \omega_1(t) \times (\mathbf{x} - \mathbf{G}_1(t))$$

Change of geometry

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Local velocity around the solid body :

$$\mathbf{u}_1(t, \mathbf{x}) = \nabla^\perp \left[\omega_1(t) |\mathbf{x} - \mathbf{G}_1(t)|^2 - \mathbf{V}_1(t) \times (\mathbf{x} - \mathbf{G}_1(t)) \right].$$

Change of geometry

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$t \mapsto (\mathbf{V}_1, \omega_1)$ is given

Local velocity around the solid body : $r = 2 \text{ diam}(B_1)$

$$(L) \quad \Lambda_1(t, \mathbf{x}) = \nabla^\perp \left[\chi_r(\mathbf{x} - \mathbf{G}_1^0) \left(\omega_1(t) |\mathbf{x} - \mathbf{G}_1(t)|^2 - \mathbf{V}_1(t) \times (\mathbf{x} - \mathbf{G}_1(t)) \right) \right].$$

where

$$\chi_r(\mathbf{y}) = 1, \quad \forall |\mathbf{y}| < r, \quad \chi_r(\mathbf{y}) = 0, \quad \forall |\mathbf{y}| > 2r.$$

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Local displacement around the solid body : X solution to

$$(X) \quad \begin{cases} X(0, \mathbf{y}) &= \mathbf{y}, \\ \dot{X}(t, \mathbf{y}) &= \Lambda_1(t, X(t, \mathbf{y})). \end{cases} \quad \forall \mathbf{y} \in \mathbb{R}^2.$$

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Remarks :

- $X(t, \cdot)$ is a smooth diffeomorphism
- $X(t, \mathbf{y}) = \mathbf{y}$, $\forall |\mathbf{y} - \mathbf{G}_1^0| > 2r$
- $X = \mathbb{I}_2 + O(|\mathbf{V}_1| + |\omega_1|)$,

Change of geometry

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Remarks :

- $X(t, \cdot)$ is a smooth diffeomorphism
- $X(t, \mathbf{y}) = \mathbf{G}_1(t) + Q_1(t)\mathbf{y}$, $\forall |\mathbf{y} - \mathbf{G}_1^0| < r - \|\mathbf{V}_1; L^1(0, t)\|$
- $X = \mathbb{I}_2 + O(|\mathbf{V}_1| + |\omega_1|)$,

Change of unknown

Given $(\mathbf{u}, p, \mathbf{G}_1, \mathbf{V}_1, \omega_1)$ such that

$$(\mathbf{V}_1, \omega_1) \in \mathcal{C}([0, T]; \mathbb{R}^2) \times \mathcal{C}([0, T])$$

we introduce X and $Y = X^{-1}$ together with:

$$\begin{aligned} \bar{\mathbf{u}}(t, \mathbf{y}) &= \nabla Y(t, X(t, \mathbf{y}))\mathbf{u}(t, X(t, \mathbf{y})), & \bar{p}(t, \mathbf{y}) &= p(t, X(t, \mathbf{y})), \\ \bar{\mathbf{V}}_1(t) &= \mathbf{Q}_1^\top(t)\mathbf{V}_1(t), & \bar{\omega}_1(t) &= \omega_1(t). \end{aligned}$$

Definition

A collection $(\mathbf{u}, p, \mathbf{G}_1, \mathbf{V}_1, \omega_1)$ is a strong solution to (v-FSIS) on $(0, T)$ if

$$(\mathbf{V}_1, \omega_1) \in \mathbf{H}^1(0, T) \times H^1(0, T), \quad \mathbf{G}_1(t) = \int_0^t \mathbf{V}_1(s) ds, \quad r > \|\mathbf{V}_1; L^1(0, T)\|,$$

if $\bar{\mathbf{u}} \in \mathcal{C}([0, T]; \mathbf{H}^1(\mathcal{F}(0))) \cap L^2(0, T; \mathbf{H}^2(\mathcal{F}(0)))$, $\bar{p} \in L^2(0, T; H_{loc}^1(\mathcal{F}(0)))$
and $(\bar{\mathbf{u}}, \bar{p}, \bar{\mathbf{V}}_1, \bar{\omega}_1)$ satisfies the following system :

Problem in fixed geometry

$$\begin{cases} \partial_t \bar{\mathbf{u}} + [\mathbf{M}\bar{\mathbf{u}}] + [\mathbf{N}\bar{\mathbf{u}}] &= [\mathbf{L}\bar{\mathbf{u}}] - [\mathbf{G}\bar{\rho}], \\ \operatorname{div} \bar{\mathbf{u}} &= 0, \end{cases} \quad \text{in } \mathbb{R}^2 \setminus \bar{B}_1,$$

$$\bar{\mathbf{u}}(t, \mathbf{y}) = \bar{\mathbf{V}}_1(t) + \bar{\omega}_1(t) \times \mathbf{y}, \quad \forall \mathbf{y} \in \partial B_1, \quad \lim_{|\mathbf{y}| \rightarrow \infty} \bar{\mathbf{u}}(t, \mathbf{y}) = 0.$$

$$\begin{cases} m_1 \frac{d}{dt} \bar{\mathbf{V}}_1 &= - \int_{\partial B_1} \Sigma(\bar{\mathbf{u}}, \bar{\rho}) \mathbf{n}_1 d\sigma_1 - m_1 \bar{\omega}_1 \times \bar{\mathbf{V}}_1, \\ J_1 \frac{d}{dt} \bar{\omega}_1 &= - \int_{\partial B_1} \mathbf{y} \times \Sigma(\bar{\mathbf{u}}, \bar{\rho}) \mathbf{n}_1 d\sigma_1, \end{cases}$$

$$\bar{\mathbf{V}}_1(0) = \mathbf{v}_1^0, \quad \bar{\omega}_1(0) = \omega_1^0, \quad \bar{\mathbf{u}}(0, \cdot) = \mathbf{u}^0,$$

where :

$$\begin{aligned} [\mathbf{M}\bar{\mathbf{u}}] &\sim \mathbf{u} \cdot \nabla \mathbf{u}, & [\mathbf{N}\bar{\mathbf{u}}] &\sim \partial_t \mathbf{u} - \partial_t \bar{\mathbf{u}}, \\ [\mathbf{L}\bar{\mathbf{u}}] &\sim \mu \Delta \mathbf{u} & [\mathbf{G}\bar{\rho}] &\sim \nabla p. \end{aligned}$$

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$$\begin{cases} \partial_t \bar{\mathbf{u}} + [\mathbf{M}\bar{\mathbf{u}}] + [\mathbf{N}\bar{\mathbf{u}}] &= [\mathbf{L}\bar{\mathbf{u}}] - [\mathbf{G}\bar{p}], \\ \operatorname{div} \bar{\mathbf{u}} &= 0, \end{cases} \quad \text{in } \mathbb{R}^2 \setminus \bar{B}_1,$$

$$\bar{\mathbf{u}}(t, \mathbf{y}) = \bar{\mathbf{V}}_1(t) + \bar{\omega}_1(t) \times \mathbf{y}, \quad \forall \mathbf{y} \in \partial B_1, \quad \lim_{|\mathbf{y}| \rightarrow \infty} \bar{\mathbf{u}}(t, \mathbf{y}) = 0.$$

$$\begin{cases} m_1 \frac{d}{dt} \bar{\mathbf{V}}_1 &= - \int_{\partial B_1} \Sigma(\bar{\mathbf{u}}, \bar{p}) \mathbf{n}_1 d\sigma_1 - m_1 \bar{\omega}_1 \times \bar{\mathbf{V}}_1, \\ J_1 \frac{d}{dt} \bar{\omega}_1 &= - \int_{\partial B_1} \mathbf{y} \times \Sigma(\bar{\mathbf{u}}, \bar{p}) \mathbf{n}_1 d\sigma_1, \end{cases}$$

$$\bar{\mathbf{V}}_1(0) = \mathbf{V}_1^0, \quad \bar{\omega}_1(0) = \omega_1^0, \quad \bar{\mathbf{u}}(0, \cdot) = \mathbf{u}^0,$$

where :

$$\begin{aligned} [\mathbf{M}\bar{\mathbf{u}}] &= o([\bar{\mathbf{u}}, \bar{\mathbf{V}}_1, \bar{\omega}_1]), & [\mathbf{N}\bar{\mathbf{u}}] &= o([\bar{\mathbf{u}}, \bar{\mathbf{V}}_1, \bar{\omega}_1]), \\ [\mathbf{L}\bar{\mathbf{u}}] &= \mu \Delta \bar{\mathbf{u}} + o([\bar{\mathbf{u}}, \bar{\mathbf{V}}_1, \bar{\omega}_1]) & [\mathbf{G}\bar{p}] &= \nabla \bar{p} + o([\bar{\mathbf{u}}, \bar{\mathbf{V}}_1, \bar{\omega}_1]). \end{aligned}$$

Problem in fixed geometry

$$\begin{cases} \partial_t \bar{\mathbf{u}} &= \mu \Delta \bar{\mathbf{u}} - \nabla \bar{p} + \mathbf{F}[\bar{\mathbf{u}}, \bar{p}, \bar{\mathbf{V}}_1, \bar{\omega}_1], \\ \operatorname{div} \bar{\mathbf{u}} &= 0, \end{cases} \quad \text{in } \mathbb{R}^2 \setminus B_1,$$

$$\bar{\mathbf{u}}(t, \mathbf{y}) = \bar{\mathbf{V}}_1(t) + \bar{\omega}_1(t) \times \mathbf{y}, \quad \forall \mathbf{y} \in \partial B_1, \quad \lim_{|\mathbf{y}| \rightarrow \infty} \bar{\mathbf{u}}(t, \mathbf{y}) = 0.$$

$$\begin{cases} m_1 \frac{d}{dt} \bar{\mathbf{V}}_1 &= - \int_{\partial B_1} \Sigma(\bar{\mathbf{u}}, \bar{p}) \mathbf{n}_1 d\sigma_1 + \mathbf{F}_M[\bar{\mathbf{V}}_1, \bar{\omega}_1], \\ J_1 \frac{d}{dt} \bar{\omega}_1 &= - \int_{\partial B_1} \mathbf{y} \times \Sigma(\bar{\mathbf{u}}, \bar{p}) \mathbf{n}_1 d\sigma_1 + T_M[\bar{\mathbf{V}}_1, \bar{\omega}_1]. \end{cases}$$

$$\bar{\mathbf{V}}_1(0) = \mathbf{V}_1^0, \quad \bar{\omega}_1(0) = \omega_1^0, \quad \bar{\mathbf{u}}(0, \cdot) = \mathbf{u}^0,$$

where :

$$\begin{aligned} [\mathbf{M}\bar{\mathbf{u}}] &= o([\bar{\mathbf{u}}, \bar{\mathbf{V}}_1, \bar{\omega}_1]), & [\mathbf{N}\bar{\mathbf{u}}] &= o([\bar{\mathbf{u}}, \bar{\mathbf{V}}_1, \bar{\omega}_1]), \\ [\mathbf{L}\bar{\mathbf{u}}] &= \mu \Delta \bar{\mathbf{u}} + o([\bar{\mathbf{u}}, \bar{\mathbf{V}}_1, \bar{\omega}_1]) & [\mathbf{G}\bar{p}] &= \nabla \bar{p} + o([\bar{\mathbf{u}}, \bar{\mathbf{V}}_1, \bar{\omega}_1]). \end{aligned}$$

Strategy for solving the new problem

$$\begin{cases} \partial_t \mathbf{u} &= \mu \Delta \mathbf{u} - \nabla \rho + \mathbf{F}, \\ \operatorname{div} \mathbf{u} &= 0, \end{cases} \quad \text{in } \mathbb{R}^2 \setminus B_1,$$

$$\mathbf{u}(t, \mathbf{y}) = \mathbf{V}_1(t) + \omega_1(t) \times \mathbf{y}, \quad \forall \mathbf{y} \in \partial B_1, \quad \lim_{|\mathbf{y}| \rightarrow \infty} \mathbf{u}(t, \mathbf{y}) = 0.$$

$$\begin{cases} m_1 \frac{d}{dt} \mathbf{V}_1 &= - \int_{\partial B_1} \Sigma(\mathbf{u}, \rho) \mathbf{n}_1 d\sigma_1 + \mathbf{F}_M, \\ J_1 \frac{d}{dt} \omega_1 &= - \int_{\partial B_1} \mathbf{y} \times \Sigma(\mathbf{u}, \rho) \mathbf{n}_1 d\sigma_1 + T_M, \end{cases}$$

$$\mathbf{V}_1(0) = \mathbf{V}_1^0, \quad \omega_1(0) = \omega_1^0, \quad \mathbf{u}(0, \cdot) = \mathbf{u}^0,$$

with polynomial nonlinearities :

$$\mathbf{F} \in \mathbf{L}^2((0, T) \times \mathcal{F}(0)), \quad \mathbf{F}_M \in \mathbf{L}^2(0, T), \quad T_M \in L^2(0, T),$$

and initial data satisfying the compatibility conditions :

$$\mathbf{u}^0 \in \mathbf{H}^1(\mathbb{R}^2 \setminus \overline{B}_1,) \quad \operatorname{div} \mathbf{u}^0 = 0, \quad \mathbf{u}^0(\mathbf{y}) = \mathbf{V}_1^0 + \omega_1^0 \times \mathbf{y} \quad \forall \mathbf{y} \in \partial B_1.$$

Analysis of the linear problem

Lemma (M. Tucsnak & T. Takahashi '03)

Given

$$\mathbf{F} \in \mathbf{L}^2((0, T) \times \mathbb{R}^2 \setminus \bar{B}_1), \quad \mathbf{F}_M \in \mathbf{L}^2(0, T), \quad T_M \in L^2(0, T)$$

and $(\mathbf{u}^0, \mathbf{v}^0, \omega^0) \in \mathbf{H}^1(\mathbb{R}^2 \setminus \bar{B}_1) \times \mathbb{R}^2 \times \mathbb{R}$ s.t.

$$\operatorname{div} \mathbf{u}^0 = 0, \quad \mathbf{u}^0(\mathbf{y}) = \mathbf{v}_1^0 + \omega_1^0 \times \mathbf{y}, \quad \forall \mathbf{y} \in \partial B_1,$$

there exists a unique solution $(\mathbf{u}, p, \mathbf{v}_1, \omega_1)$ to the previous system such that

$$\mathbf{u} \in C([0, T]; \mathbf{H}^1(\mathbb{R}^2 \setminus \bar{B}_1)) \cap L^2(0, T; \mathbf{H}^2(\mathbb{R}^2 \setminus \bar{B}_1)) \cap H^1(0, T; \mathbf{L}^2(\mathbb{R}^2 \setminus \bar{B}_1)),$$

$$p \in L^2(0, T; H_{loc}^1(\mathbb{R}^2 \setminus \bar{B}_1)),$$

$$\mathbf{v}_1 \in H^1(0, T),$$

$$\omega_1 \in H^1(0, T).$$

Moreover the mapping $(\mathbf{u}^0, \mathbf{v}^0, \omega^0, \mathbf{F}, \mathbf{F}_M, T_M) \mapsto (\mathbf{u}, \mathbf{v}_1, \omega_1)$ is linear continuous.

Introduction of a suitable operator

Functions spaces

$$\mathcal{H} = \{ \mathbf{v} \in \mathbf{L}^2(\Omega), \operatorname{div} \mathbf{v} = 0 \text{ in } \Omega, \exists (\mathbf{V}_v, \omega_v) \in \mathbb{R}^2 \times \mathbb{R}, \mathbf{v}(\mathbf{y}) = \mathbf{V}_v + \omega_v \times \mathbf{y}, \forall \mathbf{y} \in B_1 \}.$$

$$\mathcal{V} = \{ \mathbf{v} \in \mathbf{H}^1(\Omega), \operatorname{div} \mathbf{v} = 0 \text{ in } \Omega, \exists (\mathbf{V}_v, \omega_v) \in \mathbb{R}^2 \times \mathbb{R}, \mathbf{v}(\mathbf{y}) = \mathbf{V}_v + \omega_v \times \mathbf{y}, \forall \mathbf{y} \in B_1 \}.$$

Remark \mathcal{H} is a closed subspace (projector \mathbb{P}) of $\mathbf{L}^2(\mathbb{R}^2)$ with the scalar product :

$$\begin{aligned} (\mathbf{v}, \mathbf{w}) &= \int_{\mathbb{R}^2 \setminus \bar{B}_1} \mathbf{u} \cdot \mathbf{w} + \rho_1 \int_{B_1} \mathbf{v} \cdot \mathbf{w}, \\ &= \int_{\mathbb{R}^2 \setminus \bar{B}_1} \mathbf{u} \cdot \mathbf{w} + m_1 \mathbf{V}_v \cdot \mathbf{V}_w + J \omega_v \omega_w. \end{aligned}$$

Operator definition $D(A) = \{ \mathbf{v} \in \mathcal{V}, \mathbf{v} \in \mathbf{H}^2(\mathbb{R}^2 \setminus \bar{B}_1) \}$, $A = \mathbb{P}A$:

$$A\mathbf{v} = \begin{cases} \frac{\mu}{m_1} \int_{\partial B_1} 2D(\mathbf{u})\mathbf{n}_1 d\sigma_1 + \left(\frac{\mu}{J_1} \int_{\partial B_1} \mathbf{z} \times (2D(\mathbf{u})\mathbf{n}_1) d\sigma_1(\mathbf{z}) \right) \times \mathbf{y}, & \text{if } \mathbf{y} \in B_1, \\ -\mu \Delta \mathbf{u}, & \text{else.} \end{cases}$$

Analysis of A

Proposition

A is a maximal self-adjoint accretive operator.

Proof.

A is symmetric and positive. Let $(\mathbf{v}, \mathbf{w}) \in D(A)^2$ then :

$$\begin{aligned} (A\mathbf{v}, \mathbf{w}) &= -\mu \int_{\mathbb{R}^2 \setminus \bar{B}_1} \Delta \mathbf{v} \cdot \mathbf{w} + \mu \int_{\partial B_1} D(\mathbf{v}) \mathbf{n}_1 d\sigma_1 \cdot \mathbf{V}_\mathbf{w} + \mu \omega_\mathbf{w} \int_{\partial B_1} [\mathbf{z} \times D(\mathbf{v}) \mathbf{n}_1] d\sigma_1 \\ &= 2\mu \int_{\mathbb{R}^2 \setminus \bar{B}_1} D(\mathbf{v}) : D(\mathbf{w}) = \mu \int_{\mathbb{R}^2} \nabla \mathbf{v} : \nabla \mathbf{w}. \end{aligned}$$

A is maximal :

$$\forall \mathbf{f} \in \mathcal{H} \quad \exists \mathbf{v} \in D(A) \quad \mathbf{v} + A\mathbf{v} = \mathbf{f}, \text{ moreover } \|\mathbf{v}; \mathbf{H}^2(\mathbb{R}^2 \setminus \bar{B}_1)\| \leq \|\mathbf{f}; \mathcal{H}\|.$$

- **Existence and uniqueness** : variational formulation + Lax Milgram theorem
- **Regularity** : Ellipticity estimates for Stokes system



Proof of Lemma [M. Tucsnak & T. Takahashi '03]

Let $(\mathbf{F}, \mathbf{F}_M, T_M)$ source term and $(\mathbf{u}^0, \mathbf{V}^0, \omega^0)$ initial data. Set $\tilde{\mathbf{f}}$ and $\tilde{\mathbf{u}}^0$

$$\tilde{\mathbf{f}}(t, \mathbf{y}) = \begin{cases} \frac{1}{m_1} \mathbf{F}_M(t) + \frac{1}{J_1} T_M(t) \times \mathbf{y} & \text{in } B_1 \\ \mathbf{F}(t, \mathbf{y}) & \text{else} \end{cases} \quad \tilde{\mathbf{u}}^0(\mathbf{y}) = \begin{cases} \mathbf{V}^0 + \omega^0 \times \mathbf{y} & \text{in } B_1 \\ \mathbf{u}^0(\mathbf{y}) & \text{else} \end{cases}$$

Construction of the candidate solution.

There exists a unique solution to $u' = Au + \mathbb{P}f$. such that:

$$u \in \mathcal{C}([0, T]; \mathcal{V}) \cap L^2(0, T; D(A)) \cap H^1(0, T; \mathcal{H})$$

Set $\mathbf{u} = u|_{\mathbb{R}^2 \setminus \bar{B}_1}$ $\mathbf{V}_1 = \mathbf{V}_u$ $\omega_1 = \omega_u$.

$$\begin{aligned} \mathbf{u} &\in \mathcal{C}([0, T]; \mathbf{H}^1(\mathbb{R}^2 \setminus \bar{B}_1)) \cap H^1(0, T; \mathbf{L}^2(\mathbb{R}^2 \setminus \bar{B}_1)), \\ \mathbf{V}_1 &\in H^1(0, T), \\ \omega_1 &\in H^1(0, T). \end{aligned}$$

+ estimates.

Proof of Lemma [M. Tucsnak & T. Takahashi '03]

Verifications.

$\implies \operatorname{div} \mathbf{u} = 0$ and boundary conditions O.K.

Given $\mathbf{w} \in \mathcal{C}_c^\infty(\mathbb{R}^2 \setminus \bar{B}_1)$ s.t. $\operatorname{div} \mathbf{w} = 0$:

$$(u', \mathbf{w}) = - (Au, \mathbf{w}) + (\tilde{f}, \mathbf{w})$$

Proof of Lemma [M. Tucsnak & T. Takahashi '03]

Verifications.

$\implies \operatorname{div} \mathbf{u} = 0$ and boundary conditions O.K.

Given $\mathbf{w} \in \mathcal{C}_c^\infty(\mathbb{R}^2 \setminus \bar{B}_1)$ s.t. $\operatorname{div} \mathbf{w} = 0$:

$$\int_{\mathbb{R}^2 \setminus \bar{B}_1} \partial_t \mathbf{u} \cdot \mathbf{w} = - (Au, \mathbf{w}) + (\tilde{f}, \mathbf{w})$$

Proof of Lemma [M. Tucsnak & T. Takahashi '03]

Verifications.

$\implies \operatorname{div} \mathbf{u} = 0$ and boundary conditions O.K.

Given $\mathbf{w} \in \mathcal{C}_c^\infty(\mathbb{R}^2 \setminus \overline{B}_1)$ s.t. $\operatorname{div} \mathbf{w} = 0$:

$$\int_{\mathbb{R}^2 \setminus \overline{B}_1} \partial_t \mathbf{u} \cdot \mathbf{w} = \int_{\mathbb{R}^2 \setminus \overline{B}_1} \mu \Delta \mathbf{u} \cdot \mathbf{w} + (\tilde{f}, \mathbf{w})$$

Proof of Lemma [M. Tucsnak & T. Takahashi '03]

Verifications.

$\implies \operatorname{div} \mathbf{u} = 0$ and boundary conditions O.K.

Given $\mathbf{w} \in C_c^\infty(\mathbb{R}^2 \setminus \bar{B}_1)$ s.t. $\operatorname{div} \mathbf{w} = 0$:

$$\int_{\mathbb{R}^2 \setminus \bar{B}_1} \partial_t \mathbf{u} \cdot \mathbf{w} = \int_{\mathbb{R}^2 \setminus \bar{B}_1} \mu \Delta \mathbf{u} \cdot \mathbf{w} + \int_{\mathbb{R}^2 \setminus \bar{B}_1} \mathbf{f} \cdot \mathbf{w}$$

\implies there exists a pressure p s.t. $\partial_t \mathbf{u} = \mu \Delta \mathbf{u} - \nabla p + \mathbf{f}$
+ H^2 estimate

Proof of Lemma [M. Tucsnak & T. Takahashi '03]

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+ H^2 estimate

Given $\mathbf{w} \in C_c^\infty(\mathbb{R}^2)$ s.t. $\operatorname{div} \mathbf{w} = 0$ and $\mathbf{w} = \mathbf{V}$ on B_1 : $(u', \mathbf{w}) + (Au, \mathbf{w}) = (\tilde{f}, \mathbf{w})$

$$l_t + m_1 \dot{\mathbf{V}}_1 \cdot \mathbf{V} = l_\Delta - 2\mu \int_{\partial B} D(\mathbf{u}) \mathbf{n}_1 d\sigma_1 \cdot \mathbf{V} + l_f + \mathbf{F}_M \cdot \mathbf{V}$$

Proof of Lemma [M. Tucsnak & T. Takahashi '03]

Verifications.

$\implies \operatorname{div} \mathbf{u} = 0$ and boundary conditions O.K.

Given $\mathbf{w} \in C_c^\infty(\mathbb{R}^2 \setminus \bar{B}_1)$ s.t. $\operatorname{div} \mathbf{w} = 0$:

$$\int_{\mathbb{R}^2 \setminus \bar{B}_1} \partial_t \mathbf{u} \cdot \mathbf{w} = \int_{\mathbb{R}^2 \setminus \bar{B}_1} \mu \Delta \mathbf{u} \cdot \mathbf{w} + \int_{\mathbb{R}^2 \setminus \bar{B}_1} \mathbf{f} \cdot \mathbf{w}$$

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Given $\mathbf{w} \in C_c^\infty(\mathbb{R}^2)$ s.t. $\operatorname{div} \mathbf{w} = 0$ and $\mathbf{w} = \mathbf{V}$ on B_1 : $(u', \mathbf{w}) + (Au, \mathbf{w}) = (\tilde{f}, \mathbf{w})$

$$m_1 \dot{\mathbf{V}}_1 \cdot \mathbf{V} = \int_{B_1} \nabla p \cdot \mathbf{w} - 2\mu \int_{\partial B} D(\mathbf{u}) \mathbf{n}_1 d\sigma_1 \cdot \mathbf{V} + \mathbf{F}_M \cdot \mathbf{V}$$

Proof of Lemma [M. Tucsnak & T. Takahashi '03]

Verifications.

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Given $\mathbf{w} \in C_c^\infty(\mathbb{R}^2 \setminus \overline{B_1})$ s.t. $\operatorname{div} \mathbf{w} = 0$:

$$\int_{\mathbb{R}^2 \setminus \overline{B_1}} \partial_t \mathbf{u} \cdot \mathbf{w} = \int_{\mathbb{R}^2 \setminus \overline{B_1}} \mu \Delta \mathbf{u} \cdot \mathbf{w} + \int_{\mathbb{R}^2 \setminus \overline{B_1}} \mathbf{f} \cdot \mathbf{w}$$

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+ H^2 estimate

Given $\mathbf{w} \in C_c^\infty(\mathbb{R}^2)$ s.t. $\operatorname{div} \mathbf{w} = 0$ and $\mathbf{w} = \mathbf{V}$ on B_1 : $(u', \mathbf{w}) + (Au, \mathbf{w}) = (\tilde{f}, \mathbf{w})$

$$m_1 \dot{\mathbf{V}}_1 \cdot \mathbf{V} = - \int_{\partial B_1} \Sigma(\mathbf{u}, p) \mathbf{n}_1 d\sigma_1 \cdot \mathbf{V} + \mathbf{F}_M \cdot \mathbf{V}$$

Proof of Lemma [M. Tucsnak & T. Takahashi '03]

Verifications.

$\implies \operatorname{div} \mathbf{u} = 0$ and boundary conditions O.K.

Given $\mathbf{w} \in C_c^\infty(\mathbb{R}^2 \setminus \bar{B}_1)$ s.t. $\operatorname{div} \mathbf{w} = 0$:

$$\int_{\mathbb{R}^2 \setminus \bar{B}_1} \partial_t \mathbf{u} \cdot \mathbf{w} = \int_{\mathbb{R}^2 \setminus \bar{B}_1} \mu \Delta \mathbf{u} \cdot \mathbf{w} + \int_{\mathbb{R}^2 \setminus \bar{B}_1} \mathbf{f} \cdot \mathbf{w}$$

\implies there exists a pressure p s.t. $\partial_t \mathbf{u} = \mu \Delta \mathbf{u} - \nabla p + \mathbf{f}$
+ H^2 estimate

Given $\mathbf{w} \in C_c^\infty(\mathbb{R}^2)$ s.t. $\operatorname{div} \mathbf{w} = 0$ and $\mathbf{w} = \mathbf{V}$ on B_1 : $(u', \mathbf{w}) + (Au, \mathbf{w}) = (\tilde{f}, \mathbf{w})$

$$m_1 \dot{\mathbf{V}}_1 \cdot \mathbf{V} = - \int_{\partial B_1} \Sigma(\mathbf{u}, p) \mathbf{n}_1 d\sigma_1 \cdot \mathbf{V} + \mathbf{F}_M \cdot \mathbf{V}$$

\implies Newton laws.

Local existence result

Theorem ($N = 1, \Omega = \mathbb{R}^2$)

Given initial data $(\mathbf{u}^0, (\mathbf{G}_1^0, \mathbf{V}_1^0, \omega_1^0)) \in \mathbf{H}^1(\mathcal{F}(0)) \times \mathbb{R}^2 \times \mathbb{R}^2 \times \mathbb{R}$ satisfying

$$\operatorname{div} \mathbf{u}^0 = 0 \quad \forall \mathbf{x} \in \mathcal{F}(0), \quad \mathbf{u}^0(\mathbf{x}) = \mathbf{V}_1^0 + \omega_1^0 \times (\mathbf{x} - \mathbf{G}_1^0) \quad \forall \mathbf{x} \in \partial B_1(0),$$

there exists $T_0 > 0$ depending only on the size of initial data such that there exists a unique strong solution to (v-FSIS) on $(0, T_0)$ with the given initial data.

Local existence result

Theorem ($N = 1, \Omega = \mathbb{R}^2$)

Given initial data $(\mathbf{u}^0, (\mathbf{G}_1^0, \mathbf{V}_1^0, \omega_1^0)) \in \mathbf{H}^1(\mathcal{F}(0)) \times \mathbb{R}^2 \times \mathbb{R}^2 \times \mathbb{R}$ satisfying

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there exists $T_0 > 0$ depending only on the size of initial data such that there exists a unique strong solution to (v-FSIS) on $(0, T_0)$ with the given initial data.

Theorem (general case)

Given a source term $\mathbf{f} \in L^2_{loc}(0, \infty; W^{1,\infty}(\Omega) \cap L^2(\Omega))$, we assume initial data $(\mathbf{u}^0, (\mathbf{G}_i^0, \mathbf{V}_i^0, \omega_i^0)_{i \in \{1, \dots, N\}})$ satisfy a no-contact assumption and

$$\mathbf{u}^0 \in \mathbf{H}^1(\mathcal{F}(0)) \quad \operatorname{div} \mathbf{u}^0 = 0 \quad \forall \mathbf{x} \in \mathcal{F}(0),$$

$$\mathbf{u}^0(\mathbf{x}) = \mathbf{V}_i^0 + \omega_i^0 \times (\mathbf{x} - \mathbf{G}_i^0) \quad \forall \mathbf{x} \in \partial B_i(0) \quad \forall i \in \{1, \dots, N\}.$$

Then, there exists $T_0 > 0$ depending only on the size of initial data and the initial configuration such that there exists a unique strong solution to (v-FSIS) on $(0, T_0)$ with the given initial data.

Case $N = 1$, $\Omega = \mathbb{R}^2$

Theorem (Blow-up alternative)

Given $(\mathbf{u}^0, (\mathbf{G}_1^0, \mathbf{V}_1^0, \omega_1^0)) \in \mathbf{H}^1(\mathcal{F}(0)) \times \mathbb{R}^2 \times \mathbb{R}^2 \times \mathbb{R}$ satisfying

$$\operatorname{div} \mathbf{u}^0 = 0 \quad \forall \mathbf{x} \in \mathcal{F}(0), \quad \mathbf{u}^0(\mathbf{x}) = \mathbf{V}_1^0 + \omega_1^0 \times (\mathbf{x} - \mathbf{G}_1^0), \quad \forall \mathbf{x} \in \partial B_1(0),$$

there exists a unique maximal strong solution to (v-FSIS) with initial data $(\mathbf{u}^0, \mathbf{G}_1^0, \mathbf{V}_1^0, \omega_1^0)$. Moreover, the existence time T_* of the maximal solution satisfies the alternative :

- either $T_* = \infty$
- either $T_* < \infty$ and $\limsup_{t \rightarrow T_*} \|\mathbf{u}(t, \cdot) ; H^1(\mathbb{R}^2 \setminus \bar{B}_1(t))\| = \infty$.

First order estimate. Multiply momentum equation by \mathbf{u} .

$$\int_{\mathcal{F}(t)} \partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} = \int_{\mathcal{F}(t)} \operatorname{div} \Sigma(\mathbf{u}, p) \cdot \mathbf{u}$$

Case $N = 1$, $\Omega = \mathbb{R}^2$

Theorem (Blow-up alternative)

Given $(\mathbf{u}^0, (\mathbf{G}_1^0, \mathbf{V}_1^0, \omega_1^0)) \in \mathbf{H}^1(\mathcal{F}(0)) \times \mathbb{R}^2 \times \mathbb{R}^2 \times \mathbb{R}$ satisfying

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First order estimate. Multiply momentum equation by \mathbf{u} .

$$\int_{\mathcal{F}(t)} (\partial_t + \mathbf{u} \cdot \nabla) \frac{|\mathbf{u}|^2}{2} = \int_{\mathcal{F}(t)} \operatorname{div} \Sigma(\mathbf{u}, p) \cdot \mathbf{u}$$

Case $N = 1$, $\Omega = \mathbb{R}^2$

Theorem (Blow-up alternative)

Given $(\mathbf{u}^0, (\mathbf{G}_1^0, \mathbf{V}_1^0, \omega_1^0)) \in \mathbf{H}^1(\mathcal{F}(0)) \times \mathbb{R}^2 \times \mathbb{R}^2 \times \mathbb{R}$ satisfying

$$\operatorname{div} \mathbf{u}^0 = 0 \quad \forall \mathbf{x} \in \mathcal{F}(0), \quad \mathbf{u}^0(\mathbf{x}) = \mathbf{V}_1^0 + \omega_1^0 \times (\mathbf{x} - \mathbf{G}_1^0), \quad \forall \mathbf{x} \in \partial B_1(0),$$

there exists a unique maximal strong solution to (v-FSIS) with initial data $(\mathbf{u}^0, \mathbf{G}_1^0, \mathbf{V}_1^0, \omega_1^0)$. Moreover, the existence time T_* of the maximal solution satisfies the alternative :

- either $T_* = \infty$
- either $T_* < \infty$ and $\limsup_{t \rightarrow T_*} \|\mathbf{u}(t, \cdot) ; H^1(\mathbb{R}^2 \setminus \bar{B}_1(t))\| = \infty$.

First order estimate. Multiply momentum equation by \mathbf{u} .

$$\frac{1}{2} \frac{d}{dt} \left[\int_{\mathcal{F}(t)} |\mathbf{u}|^2 \right] = \int_{\mathcal{F}(t)} \operatorname{div} \Sigma(\mathbf{u}, p) \cdot \mathbf{u}$$

Case $N = 1$, $\Omega = \mathbb{R}^2$

Theorem (Blow-up alternative)

Given $(\mathbf{u}^0, (\mathbf{G}_1^0, \mathbf{V}_1^0, \omega_1^0)) \in \mathbf{H}^1(\mathcal{F}(0)) \times \mathbb{R}^2 \times \mathbb{R}^2 \times \mathbb{R}$ satisfying

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there exists a unique maximal strong solution to (v-FSIS) with initial data $(\mathbf{u}^0, \mathbf{G}_1^0, \mathbf{V}_1^0, \omega_1^0)$. Moreover, the existence time T_* of the maximal solution satisfies the alternative :

- either $T_* = \infty$
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Formal estimates

$$B_1 = B(0, R)$$

Assume $(\mathbf{u}, p, \mathbf{G}_1, \mathbf{V}_1, \omega_1)$ is a solution to (v-FSIS)

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$$\sup_{(0, T)} \|\mathbf{u}^0; H^1(\mathcal{F}(t))\| \leq C(\mathbf{u}_0, \mathbf{V}_1^0, \omega_1^0, T).$$

Case $N \geq 1$ and $\Omega \subset \subset \mathbb{R}^2$

With a similar construction, it yields :

Theorem

Given a source term $\mathbf{f} \in L^2_{loc}(0, \infty; W^{1, \infty}(\Omega) \cap L^2(\Omega))$, we assume initial data $(\mathbf{u}^0, (\mathbf{G}_i^0, \mathbf{V}_i^0, \omega_i^0)_{i \in \{1, \dots, N\}})$ satisfying a no-contact assumption and

$$\mathbf{u}^0 \in \mathbf{H}^1(\mathcal{F}(0)), \quad \operatorname{div} \mathbf{u}^0 = 0, \quad \forall \mathbf{x} \in \mathcal{F}(0),$$

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Then, there exists a maximal T_* such that there exists a unique strong solution to (v-FSIS) on $(0, T_*)$ with the given initial data. Moreover, there holds the alternative

- either $T_* = +\infty$,
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$$\frac{1}{2} \frac{d}{dt} \left[\sum_{i=1}^N \left(m_i |\mathbf{V}_i|^2 + J_i |\omega_i|^2 \right) + \int_{\mathcal{F}(t)} |\mathbf{u}|^2 \right] + 2 \sum_{i=1}^N |B_i| \omega_i^2 + \mu \int_{\mathcal{F}(t)} |\nabla \mathbf{u}|^2 = 0.$$

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Question : Does contact mean a "real" blow-up ?